



The bridge to design

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- Multiscale approaches for nanostructures
 - Simulation of nanostructures
 - Cross Bar Architectures
 - Simulation of organic FETs and organic inverters
- Field coupling in Quantum Cellular Automata
 - o Quantum Cellular Automata
 - Magnetic logic gate
- Conclusions



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International Technology Roadmap for Semiconductors

ТЛП



Time

2002

ISS US

P.Gargini





Beyond CMOS logic and memory device candidates:

- Nanowire transistors
- CNT transistors
- Resonant tunneling devices
- NEMS devices
- Single electron transistors
- Molecular devices
- Spintronic devices

All those candidates (some of which not yet demonstrated) still suffer from major reliability and stability problems





- Circuit-based modeling (SPICE)
- Drift-diffusion device simulation (OTFT, MOSFET, HEMT,..)
- □ Monte Carlo device simulation (HEMT, QCL, …)
- Ab-initio and quasi-ab-initio modeling of molecules
- NEGF calculation of carrier transport in molecules
- Tight binding calculation of electronic properties of nanostructures







An electronic structure problem

Self-consistent solution

Efficient Density Functional Theory codes

SIESTA (Soler et al.) \rightarrow Numerical basis functions, O(n) scaling

DFTB (Frauenheim et. al.) →DFT parameterized tight-binding









The simulated structure is divided into 'bulk' and 'device regions'



Self-consistent potential of molecule contacted to carbon nanotubes



Transmission function





The current-voltage characteristics can be calculated from the so called Landauer equation as an integral of the transmission function T(E)







Al long as no resonances are included in the transmission window (low bias), the current is negligible (*tunneling regime*). The threshold voltage occurs when the HOMO-LUMO states enter such window. The terrace is due to the fact that the transmission peaks broaden and shrink at increasing voltage



-4.95

-5.05

-5.1 [A] AG-5.15 -5.2 -5.1

-5.2

-5.25 -5.3

Cf source

-5

h

 $\mathbf{E}_{\mathrm{f,source}}$

60

70

50

40

The CNTs behave like ideal p-type nanowires. The gate voltage modulates the transmission of the holes which are injected from the source contact.



The energie-profile of the highest valence-subband in a (10,0) CNT for $V_{DS} = 0.2$ V and $V_{GS} = 0.6$ V.

30 z [ang]

aate

20







Landauer Formalism is used to calculate the current between souce and drain







We can simulate transport in molecular structures quite reliably, accounting for charging effects, dissipation, screening as long as the molecular systems consist of a sufficiently small number of atoms,



the simulation of realistic systems, especially in terms of the contacts and in general the coupling with the external environment is still a problem



Multiscale modeling











We use SPICE circuit simulations to investigate such circuits. We built the model of the junction using published (molecular) I-V curves and then generated SPICE netlists of large circuits using a Matlab script.









Scheme I

- Overall leakage can be high
- Parasitic currents does not change the readcurrent significantly
- Figure of merit for diode: I_{on}/I_{off}

Scheme II

- Overall leakage is very low
- Parasitic currents limit scalability
- Figure of merit for diode: $I_{on}/I_{on/2}$

Scheme III

- Overall leakage is very low
 - Parasitic currents have little
 effect
- Capacitive effects may play
 an important role



Requirements





- The junction should deliver high currents at low voltage, ideally below 6 V.
- Stability with time (10 years) and to electrical cycling (10⁶⁻⁸ for reading, 10⁴ for writing)



Bit lines

We have set up a circuit model for molecular cross bar arrays to address the performance and the scalability of such systems. The IV characteristics for the SPICE model taken from: E. Lortscher, J. W. Ciszek, J. Tour, H. Riel: Reversible and Controllable Switching of a Single-Molecule Junction, Small 2, No. 8-9, 973 (2006)



- Nonlinearity of the molecule IV characteristics enhances read-out margin
- Coupling molecular layer with solid-state diodes enables large-scale circuits



ПΠ

With all the information gathered it is possible to optimized the diode and know what are the characteristics that it must have.

- Low contact resistance (<10⁻⁵ Ωcm²)*
- Carrier concentration $< 10^{17}$ cm⁻³).
- o High Schottky barrier (ZnO/Ag \rightarrow 0.9 eV).

*<u>Carrier concentration dependence of</u> <u>*Ti/Al/Pt/Au contact resistance on n-type ZnO* K. Ip, Y. Heo, K. Baik, D. P. Norton, S. J. Pearton, and F. Ren, Appl. Phys. Lett. **84**, 2004.</u>



We implemented a Matlab-script, which is able to generate circuits up to a million nodes (junctions). Behavior of large crossbar arrays can be analyzed this way.









For circuits larger than ca. 100×100 junctions, the difference between 0 and 1 written bit value disappears – in a worst-case scenario, parasitic currents overwhelm the current to be sensed.



For **optimized** diodes, a 1 Mbit array seems plausible with any bias scheme











The simulation employs a modified *drift-diffusion model* implemented in the industry tool CENTAURUS (former ISETCAD).

Input parameters are

device geometry,

 Q_I

- electronic levels in Pentacene, and
- properties of the material interfaces.



SiO₂/Pentacene interface

fixed interface charges $p_{At} = N_{At} f_n$ acceptor-type traps, i.e. negative if filled, neutral if empty



The Drift-Diffusion Approach



$$J_{n} = -q\mu_{n}\nabla V + qD_{n}\nabla n$$

$$J_{p} = -q\mu_{p}\nabla V - qD_{p}\nabla p$$

$$\Delta\varepsilon_{0}\varepsilon_{r}V = -q(p-n-p_{At}+Q_{I})$$

$$\nabla \cdot J_{n} = q(R-G)$$

$$\nabla \cdot J_{p} = -q(R-G)$$

The self-consistent simulator scheme numerically solves the coupled equations on a spatial grid which models the device geometry



A. Bolognesi, et al. J. Comp. Elec. 2, 297 (2003)

- C. Erlen, et al. J. Comp. Elec. 5, 345 (2006)
- P. Lugli, et al. Appl. Phys. A, Vol. 87, No. 3., 2007





Effective DOS = $1e21 \text{ cm}^{-3}$ Contact Barrier = 0.1 eV

Lee et al. APL 84(10):1701, 2004

Field-dependent mobility has been implemented according to

$$\mu = \mu_0 \cdot e^{\sqrt{\frac{E}{E_{critical}}}}$$

The low-field mobility μ_0 is a transistor property and has to be extracted via simulations.

Bolognesi et al. IEEE Trans. El. Dev., 51(12):1997, 2004







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Single electron transistors

Experimental demonstration for reduced complexity \rightarrow it works only at subkryogenic temperatures

Coupled nanoscale magnets

Cowburn et. al, Science





Very promising, experimentally demonstrated at room temperature \rightarrow slow





Idea of field-coupling













Coupled nanomagnets

ТШ



Investigations of permalloy nanomagnets (thermally evaporated and patterned by electron beam lithography) confirm the simulation results













Programmable 2-input AND or OR gate.

With the majority gate, all Boolean logic functions could be realized

Benefits of QCA architecture:

- ➢No static dissipation
- ≻No Interconnection problem

Potentially simple fabrication and scalability to molecular level





Output points down only if both inputs are pointing up \rightarrow NAND gate.

•Difficult to design – ferro- and antiferromagnetic couplings to the central dot should be equally strong

•Electrical inputs are difficult to fabricate – horizontally lying dots provide a hard-wired input. No output, we just imaged it with the MFM

 Design is based on Parish and Forshaw:
 200 nm Magnetic Cellular Automate Systems IEE Proc.-Circuits Devices Syst., Vol. 151, No. 5, October 2004

Imre et. al. Science 2006







SEM images

MFM images



Two vector fields: $\mathbf{M}(\mathbf{r},t) = \mathbf{H}_{eff}(\mathbf{r},t)$

$$\frac{\partial \mathbf{M}(\mathbf{r},t)}{\partial t} = -\gamma \mathbf{M}(\mathbf{r},t) \times \mathbf{H}_{\text{eff}}(\mathbf{r},t) - \frac{\alpha \gamma}{M_s} \Big[\mathbf{M}(\mathbf{r},t) \times \mathbf{M}(\mathbf{r},t) \times \mathbf{H}_{\text{eff}}(\mathbf{r},t) \Big]$$

Landau – Lifshitz equation: how the magnetization distribution changes under the influence of an effective field $\mathbf{H}_{\text{eff}}(\mathbf{r},t)$

The effective field:

$$\mathbf{H}_{eff}(\mathbf{r},t) = \mathbf{H}(\mathbf{r},t) + \mathbf{H}_{exch}(\mathbf{r},t)$$
Maxwell equations
"Real" magnetic field $\mathbf{H}(\mathbf{r},t) = \frac{\mathbf{B}(\mathbf{r},t)}{\mu_0} - \mathbf{M}(\mathbf{r},t)$

$$\operatorname{div} \mathbf{B}(\mathbf{r},t) = 0$$
Magnetic domain theory





Design rules:

- •Biasing dot is smaller than inputs
- •Output is of the gate is shifted to the left \rightarrow switches later





 \rightarrow

 \rightarrow

 \rightarrow

The challenges:

- How to make signals propagating?
- How to write in the magnets?
- How to read out the magnets?

- Integrated clocking
- Localized field from wires
 - Hall sensor





Rapidly moving domain walls are the main source of dissipation in magnetic materials

>Make the magnets sufficiently small (submicron size magnets has no internal domain walls)

Switch them slowly (use adiabatic pumping)





ПП

Architecture: Layer of pumping wires above the dots:



Signal propagation



Wires placed above / below the dots can provide a wave-like propagating magnetic field, which switches the magnets one after the other.

This simulation shows the horizontal magnetic field distribution above the wires.



Propagation of the Ordering Wave



Switching off the current of the wires on after the other creates a magnetic field 'wave' in the plane of the magnets.

At the wavefront magnets are switching to their up / down state – even a long wire becomes perfectly ordered.

- A variety of nano devices can be reliably fabricated from various materials
- Novel circuits and architectures are going to be needed for a full exploitation of nanocomponents
- Several open questions still exist concerning the mode of operation of such devices
- Modeling and simulation can provide important answers
- A multiscale approach is needed in order to describe realistic systems